#### **Coronal equilibrium**

- Coronal equilibrium occurs when heat or mechanical energy sets a kinetic temperature
- Besides the sun, this occurs in ...
- Accretion disks
  - Gas "viscosity" converted into heat
- Shocks
  - Collision between two parcels of gas,
  - Energy of motion converted into heat

#### Shock in an ideal gas

- $\blacklozenge$  Energy of motion  $^{1\!\!/_2}\rho\,v^2$
- Thermal energy 3/2 nkT
- T(shock) =  $\rho v^2 / 3nk = v^2 / 3k \rho/n = v^2 \mu / 3k$

• If the gas is ionized then  $n \approx 2n(H)$  so  $\mu \approx m_p / 2$ 

#### All stars have winds

- Wind speed roughly escape velocity from surface of star
- Wind strikes surrounding clouds creating shocks
- visible in x-ray for moderate to high speed shocks, in the optical / UV for slow shocks

Earth within solar wind

◆ Solar wind speed ≈400 km/s at Earth's orbit



























#### **Planetary nebula**

◆ Wind~10<sup>3</sup> km/s



## T Tauri jets • ≈ 50 km/s



#### Homework

- Choose one of these objects and compute its shock spectrum
- Plot the ionization distribution and the spectrum



#### What our calculation included

- Coronal equilibrium
- Lightest 30 elements
- All stages of ionization
- About 100 molecules
- Fully self consistent
- Time steady
- Ionization caused by thermal collisions
- Light was not considered

Mon. Not. R. astr. Soc. (1969) 142, 501-521.

THE IONIZATION EQUILIBRIUM OF ELEMENTS BETWEEN CARBON AND NICKEL

Carole Jordan

(Received 1968 July 31)

SUMMARY The ionization equilibrium N(ion)/N(Element), for ions of elements abundant in the Sun, has been calculated as a function of temperature between  $T_{\pi^*} \sim 10^4$  K and  $\sim 10^4$  K. Two sets of results are given. The first includes the processes of collisional ionization, collisional excitation followed by auto-ionization, direct radiative recombination, radiative recombination via bound levels, and di-electronic recombination at the electron temperature in the solar atmosphere enters into the calculations these results are applicable only contained in the difference recombination rate. These results are applicable to any low density plasma where the radiation field is negligible.

http://adsabs.harvard.edu/abs/1969MNRAS.142..501J

	- log N(1)/N(E) for 0 ions									
log10Te	Ion	01	ıt	ш	17	v	VI	VII	VI 11	11
3.7		0.00	-	-	-	-	-	-	-	-
3.8		0.00	6.15	-	-	-	-		-	-
3.9		0.00	3.68	-	-	-	-	- 1	-	-
4.0		0,01	1.71	-	-	-	~	-	-	-
4.1		0.20	0.44	-	-	-	-	-	-	-
4.2		1.04	0.04	-	-	-	<b>a</b> -	-	-	-
4.3		2,01	0.00	-	-	-	-	-	-	-
4.4		2.79	0.00	3.08	-	-	-	-	-	-
4.5		3,48	0.01	1.58	-	-	-			
4.6		-	0,14	0.55	4.30	-	-	-	-	-
4.7		-	0,59	0,13	2,55	-	-	-	-	-
4+8		-	1,15	0.04	1.43	5.33	-	-	-	-
4.9		-	1,78	0.13	0,60	3.35	-	-	-	-
5.0		-	2,47	0.38	0.25	1.90	-	-	-	-
5.1		-	3,29	0.83	0,12	1,04	4.13	-	-	-
5.2		-	4,17	1.39	0.21	0.46	2.51	-	-	-
5.3		-	-	2,11	0,52	0,19	1.27	2,56	-	-
5.4		-	-	3.12	1,11	0.21	0,66	1.03	-	-
5.5		-	-	4.34	2,01	0.79	0,60	0.25	-	-
5.6		-	-	5.89	3.26	1,66	1,01	0.05	-	-
5-7			-	-	4.42	2.49	1.41	0,02	-	-
5.8		-	-	-	5.56	3.32	1.85	0,00	-	-
5.9		-	-	-	-	4.03	2,23	0.00	3.22	-
6.0		-	-	-	-	4.63	2.56	0,00	2,10	5.49
6.1		-	•	-	-	-	2.78	0,03	1.21	3.61
6.2		-	-	-	-	-	-	0,12	0,63	2.12
6.3		-	-	-	-	-	-	0,34	0.32	1.13
6.4		-	-	-	-	-	-	0.74	0,29	0.50
6.5		-	-	-	-	-	-	1,28	0.51	0.19
6.6		-	-	-	-	-	-	1,89	0.82	0,08
6.7		-	-	-	-	-	-	2.47	1.14	0.03
6.8		-	-	-	-	-	-	3.04	1.46	0.02
6,9		-	-	-	-	-	-	3.55	1,75	0.01
7.0		-	-	-	-	-		4.01	2.00	0.00



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	#Te	0	0+1	0+2	0+3	0+4	0+5	0+6	0+7	0+8
	3.08	-5.95								
	3.10	-5.30								
	3.28	-2.53								
	3.30	-1.39								
	3.48	-0.84								
	2.50	-0.46								
	3.78	-0.48								
	3,88	-0,49								
	3.98	-0.40	-8.41							
	4.08	-0.89	-3.55							
	4.18	-0.02	-1.42							
	4.28	-0.26	-0.35	7.04						
	4.30	-0.99	-0.05	-/.21						
	4.58	-2.46	-0.00	-2.96		-	-	-		
	4.68	-2.93	-0.02	-1.33	-5.77					
	4.78	-3.28	-0.18	-0.47	-3.41					
	4.88	-3.70	-0.53	-0.16	-2.02	-6.65				
	4.98	-4.21	-0.97	-0.89	-1.11	-4.45				
	5.08	-4.84	-1.49	-0.18	-0.51	-2.82	-7.27			
	5.18	-5.62	-2.13	-0.45	-0.21	-1.66	-4.85	-8.35		
	5.38	-7.52	-2.8/	-1.40	-0.27	-0.00	-1.70	-2.91		
	5.48	-8.70	-4.71	-2.13	-0.64	-0.25	-8.84	-1.21		
	5.58		-6.11	-3.29	-1.47	-0.64	-8.62	-8.31		
	5.68		-7.87	-4.81	-2.71	-1.48	-0.95	-8.07	-8.37	
	5.78			-6.29	-3.93	-2.35	-1.38	-0.02	-6.26	
	5.88			-7.61	-5.02	-3.12	-1.77	-0.01	-4.59	
	5.98			-8.77	-5.95	-3.77	-2.10	-0.00	-3.23	-/./0
	6.18				-7.30	-4.71	-2.33	-0.01	-1.30	-3.56
	6.28				-7.96	-5,86	-2.60	-8.18	-8.71	-2.14
	6.38				-8.55	-5.46	-2.80	-8.28	-8.39	-1.13
	6.48					-5.99	-3.15	-0.63	-8.34	-8.51
	6.58					-6.67	-3.67	-1.13	-8.51	-8.21
	6.68					-7.41	-4.28	-1.69	-8.79	-8.09
	6.78					-8.15	-4.88	-2.24	-1.08	-8.04
	6.88					-8.85	-5.45	-2.75	-1.36	-8.02
	0.98						-3.99	-3.22	-1.61	-0.01



#### Shock in an ideal gas

- $\blacklozenge$  Energy of motion  $^{1\!\!/_2} \rho \ v^2$
- Thermal energy 3/2 nkT
- $\label{eq:theta} \black{T(shock)} = \rho \ v^2 \ / \ 3nk = v^2 \ / \ 3k \ \rho/n = \ v^2 \ \mu \ / \ 3k$
- If the gas is ionized then  $n \approx 2n(H)$  so  $\mu \approx m_p / 2$
- ◆ T(shock) = 20 v(km/s)<sup>2</sup> K

#### Watch for these correlations

•  $T_{\text{shock}} \sim = 20 \text{ v}(\text{km/s})^2 \text{ [K]}$ 

Ionization potential Z<sup>2</sup> Ryd
 ~≈1.5×10<sup>5</sup> Z<sup>2</sup> [K] (from energy match)
 ~≈10<sup>4</sup> Z<sup>2</sup> [K] (thermal distribution of energies)

◆ 100 km/s -> 200 000 [K]

◆ From 200 000 K = 10<sup>4</sup> Z<sup>2</sup> -> Z≈4

#### Watch for these correlations

•  $T_{shock} \sim = 20 \text{ v}(\text{km/s})^2 \text{ K}$ 

- E(level above ground) = Z<sup>2</sup> (1-n<sup>-2</sup>)
- Rydberg formula -E(photon) =  $E_u - E_i = Z^2 (n_i^{-2} - n_u^{-2})$
- High Z, lines have higher energy, shorter wavelength

Earth within solar wind

◆ Solar wind speed ≈400 km/s at Earth's orbit





#### T Tauri jets





# BN – KL in Orion • 10 km/s

### **Planetary nebula**

♦ Wind~10<sup>3</sup> km/s

































**Photoionization rate** [s<sup>-1</sup>]

• Flux of ionizing photons × photoionization cross section

$$\int_{\nu_0}^{\infty} \frac{4\pi J_{\nu}}{h\nu} a_{\nu}(\mathbf{H}^0) d\nu = \int_{\nu_0}^{\infty} \phi_{\nu} a_{\nu}(\mathbf{H}^0) \ d\nu = \Gamma(\mathbf{H}^0)$$
(2.1)

#### Hydrogen collisional ionization



- Electron collisional ionization rate
   -c = k<sub>8</sub> n<sub>e</sub>
- Hydrogen atom collisional rate
   -c = k<sub>9</sub> n(H<sup>0</sup>)
- Palla+1983

#### **Radiative recombination**

- Brems is produced by distant collisions between electrons and charges
  - Photon energy taken from kinetic energy of free electron
- A very close collision can cause an electron to lose so much energy that it becomes bound to the ion
- Ion + electron => atom + photon
- Rate coefficient α cm<sup>3</sup> s<sup>-1</sup>
- Rate =  $\alpha n_e$

<b>Radiative recombination</b>	
$\begin{split} \alpha_B \left( T_e \right)  =  \left\{ \begin{array}{l} 2.90  \times  10^{-10}   T_e^{ -0.77} \\ 1.31  \times  10^{-8}   T_e^{ -1.13} \end{array} \right, \\ T_e & \leq 2.6  \times  10^4  \mathrm{K} \\ T_e & > 2.6  \times  10^4  \mathrm{K} \end{array} \right\}  \mathrm{cm^3 \ s^{-1}}  . \end{split}$	(6)
<ul> <li>Ferland 1980</li> </ul>	

#### **Recombination AGN3 Chap 2**

- Electron and ion recombine, emitting energy
- Radiative recombination for H and He
- Dielectronic recombination for heavy elements Table 2.7

			T		
	1,250 K	2,500 K	5,000 K	10,000 K	20,000 K
$\alpha_A = \sum_{1}^{\infty} \alpha_n$	$1.74\times 10^{-12}$	$1.10\times 10^{-12}$	$6.82  imes 10^{-13}$	$4.18\times10^{-13}$	2.51 × 10 <sup>-13</sup>
$\alpha_B = \sum_{2}^{\infty} \alpha_n$	$1.28\times 10^{-12}$	$7.72\times10^{-13}$	$4.54\times10^{-13}$	$2.59\times 10^{-13}$	1.43 × 10 <sup>-13</sup>
$\alpha_C = \sum_{3}^{\infty} \alpha_n^{\circ}$	$1.03\times 10^{-12}$	$5.99  imes 10^{-13}$	$3.37\times10^{-13}$	$1.87  imes 10^{-13}$	9.50 × 10 <sup>-14</sup>
$\alpha_D = \sum_{4}^{\infty} \alpha_n$	$8.65\times10^{-13}$	$4.86\times10^{-13}$	$2.64  imes 10^{-13}$	$1.37 \times 10^{-13}$	6.83 × 10 <sup>-14</sup>













#### Let's model a ...

- Relatively dense,  $n_{\rm H} = 10^4 \, {\rm cm}^{-3}$
- ISM cloud
- O6 star, but all light at 2 Ryd
- ◆ Q(H)=10<sup>48</sup> s<sup>-1</sup>
- ◆ 1 pc = 3×10<sup>18</sup> cm away



#### Strömgren length



#### Strömgren length

 Number of ionizing photons entering layer is balance by number of recombinations along it

L

$$\varphi(H) = \Lambda e \cap \varphi \times$$



#### Beyond the H<sup>+</sup> layer

- Little H<sup>+</sup> ionizing radiation gets past the H<sup>+</sup> layer
- Deeper regions are atomic or molecular
- Also cold and produce little visible light
- Large extinction due to dust

